Unveiling community structures in weighted networks

Nelson A. Alves
Departamento de Física e Matemática, FFCLRP Universidade de São Paulo, Avenida Bandeirantes 3900, CEP 14040-901, Ribeirão Preto, São Paulo, Brazil

(Received 28 February 2007; published 4 September 2007)

Random walks on simple graphs in connection with electrical resistor networks lead to the definition of Markov chains with transition probability matrix in terms of electrical conductances. We extend this definition to an effective transition matrix P_{ij} to account for the probability of going from vertex i to any vertex j of the original connected graph G. Also, we present an algorithm based on the definition of this effective transition matrix among vertices in the network to extract a topological feature related to the manner by which graph G has been organized. This topological feature corresponds to the communities in the graph.

DOI: 10.1103/PhysRevE.76.036101 PACS number(s): 89.75.Hc, 05.10.−a

I. INTRODUCTION

Network modeling is becoming an essential tool to study and understand the complexity of many natural and artificial systems [1]. Applications [2–5] include technological networks such as the Internet, the World Wide Web, and the electric power grid; biological networks such as metabolic [6–8] and amino acid residue networks [9–12]; and far more studied, social networks. This understanding follows from the statistical characterization of their topological properties usually related to complex networks. In general, the statistical analysis includes the degree distribution $P(k)$, the average degree $\langle k \rangle$, the clustering coefficient C, the "betweenness" of a vertex i and "assortative mixing" describing the correlations among vertices in the network.

Currently, an important issue within complex network field is the study and identification of community structure, a problem also known as graph partitioning. Many definitions of community are presented in the literature. In essence, a community is a group of vertices that are more highly connected to each other than to vertices in other groups. A successful method to identify communities should find automatically the meaningful groups of the network. Moreover, the method needs to perform that task in computer times that are not prohibitive for large network sizes.

Various methods have been proposed to identify communities in networks [13,14]. In particular, some methods have been based on betweenness measures [15], random walks [16,17], resistor network [18], Laplacian eigenvalues [19,20], quantitative definitions of community structures in networks [21] or on benefit functions, where modularity is an example [15,19]. Those methods discover communities in computation time that typically scales with the network size n as n^3 or even n^4. There is a method that scales linearly in time but needs a parameter-dependent consideration [18]. This method views the network as a two-dimensional (2D) resistance circuit with current flowing through all edges represented by resistors. The automatic community finding procedure is hampered by the need of electing two nodes (poles) that lie in different communities and by the introduction of a threshold in the voltage spectrum.

Here we show how random walkers on graphs, also in connection with electrical networks, unveil the network community structure throughout an hierarchical connection of vertices. Our method combines the Laplacian eigenvalue approach with electrical network theory. A brief review of how the spectral graph theory can characterize the structural properties of graphs using the eigenvectors of the Laplacian matrix has been presented by Newman [19].

Our method relies on a generalization of the usual transition probability matrix **P**. The matrix element P_{ij} is the probability for a walk from vertex i on a weighted graph to an adjacent vertex j. We introduce an effective transition matrix that also accounts for hops on the graph. This is achieved by evaluating the conductances, the inverse of resistances, between any two vertices. By defining a similarity matrix as a function of the effective transition matrix elements, it is possible to extract a topological feature related to the manner graph G which has been organized. It turns out that this topological feature corresponds to hierarchically connected classes of vertices which form communities in the network.

To explain our method, we present the essential of the spectral analysis of Laplacian matrices in Sec. II. In Sec. III we present the arguments leading to the similarity matrix that sets a scale to extract the community structure. In Sec. IV we describe how to implement the algorithm and show the results for the karate club network studied by Zachary [22] and for the model designed by Ravasz and Barabási [23], an example of network with the scale-free property and modular structure. Section V concentrates our discussions on weighted graphs and Sec. VI contains our conclusions.

II. LAPLACIAN EIGENVALUES AND TRANSITION MATRIX

Let us consider a simple graph G, i.e., undirected and with no loops or multiple edges, on a finite vertex set $V=\{1,2,\ldots,n\}$ and edge set E, represented by the adjacency matrix **A**. The degree k_i for each vertex i is obtained from the adjacency matrix **A** as $k_i=\Sigma_{j=1}^{n}A_{ij}$. For nonweighted graphs, the symmetric $n\times n$ adjacency matrix takes values $A_{ij}=1$, if there is an edge connecting vertices (i,j) and 0 otherwise. Thus, k_i counts the number of edges that connect the selected vertex i to other vertices. This extends naturally to the weighted adjacency matrix which we treat in Sec. V.

For our purpose we study the graph G through a positive semidefinite matrix representation. This is achieved in the usual manner using the Laplacian. The Laplacian matrix of a graph G on n vertices, denoted by $\mathbf{L}(G)$, is simply the matrix with elements

$$L_{ij} = \begin{cases} k_i & \text{if } i = j, \\ -1 & \text{if } i \text{ and } j \text{ are adjacents}, \\ 0 & \text{otherwise}, \end{cases} \quad (1)$$

which corresponds to the degree diagonal matrix minus the adjacency matrix, $\mathbf{L}=\mathbf{K}-\mathbf{A}$. The Laplacian matrix has a long history. It was introduced by Kirchhoff in 1847 with a paper related to electrical networks [24] and consequently is also known as the Kirchhoff matrix.

The Laplacian matrix is real and symmetric. Moreover, $\mathbf{L}$ is a positive semidefinite singular matrix with n eigenvalues λ_i and eigenvectors v_i. If we label the eigenvalues in increasing order $\lambda_1 \leq \lambda_2 \leq \cdots \leq \lambda_n$, we have $\mathbf{L}(G)v_1=0$. The eigenvalue $\lambda_1=0$ is always the smallest one and has the normalized eigenvector $v_1=(1,1,\ldots,1)/\sqrt{n}$. Since the matrix $\mathbf{L}(G)$ is singular, it has no inverse, but in such cases it is possible to introduce the so-called generalized inverse ($\mathbf{L}^\dagger$) of $\mathbf{L}$ according to the Moore and Penrose definition [25].

Among many properties for the second smallest eigenvalue $\lambda_2(G)$, known as the algebraic connectivity, we recall that [24,26] $\lambda_2(G)=0$, if and only if, G is not connected. For connected networks, the eigenvector components of the first non-null eigenvalue (λ_2) have been applied as an approximate method for grouping vertices into communities [19,20,27]. However, the success in partitioning depends on how well λ_2 is separated from other eigenvalues.

From now on we identify the graph $G=(V,E)$ with an electrical network connected by edges of unit resistances [28,29]. A random walk on G is a sequence of states (vertices) chosen among their adjacent neighbors. To describe the overall behavior of a walker on G, one needs to go beyond the usual analysis of Markov chains with a transition matrix P_{ij}, probability to go from vertex i to an adjacent vertex j, to include also hops, i.e., moves across the graph. For this end, we evaluate the effective resistances r_{ij} between all distinct vertices i and j of G. Those effective resistances r_{ij} can be numerically evaluated by means of the electrical network theory as [30,31]

$$r_{ij} = (L^\dagger)_{ii} + (L^\dagger)_{jj} - (L^\dagger)_{ij} - (L^\dagger)_{ji}, \quad (2)$$

for $i \neq j$ and $r_{ij}=0$ for $i=j$. Here, $\mathbf{L}^\dagger(G)$ is the Moore-Penrose generalized inverse of the Laplacian matrix $\mathbf{L}(G)$. Its definition amounts to write $\mathbf{L}^\dagger(G)$ as

$$(L^\dagger)_{ij} = \sum_{k=1}^{n-1} \frac{1}{\lambda_k} v_{ki} v_{kj}. \quad (3)$$

This leads to a simple formulation of the effective resistances between all pairs of vertices as a function of the eigenvalues and eigenvectors of $\mathbf{L}(G)$,

$$r_{ij} = \sum_{k=1}^{n-1} \frac{1}{\lambda_k} (v_{ki} - v_{kj})^2. \quad (4)$$

As a natural generalization, it is convenient to define the effective conductances c_{ij} for all pairs of vertices (i,j) as $c_{ij}=1/r_{ij}$, for $i \neq j$.

As a consequence of the above results it is possible to extend the usual random process that moves around through adjacent states i and j to hops on the graph. We define the hop transition probability from vertex i to any vertex j by

$$P_{ij} = \frac{c_{ij}}{c_i}, \quad (5)$$

where c_{ij} is the effective conductance from i to j and $c_i = \Sigma_j c_{ij}$. Since a connected network is considered, the probability that a walker who begins the run at any given vertex i and reaches any other given vertex does not vanish.

III. METHOD

Although P_{ij} is not necessarily equal to P_{ji}, it is possible to describe hierarchical classes of states perceived by the walker as follows.

First, we consider the generalized "distance" expression,

$$d_{ij}^{(q)} = \frac{\left(\sum_{k \neq i,j}^{n} |P_{ik} - P_{jk}|^q \right)^{1/q}}{n-2}, \quad (6)$$

where q is a positive real number, as a similarity measure between any vertices. Small $d_{ij}^{(q)}$ would imply high similarity between i and j and could be used to set a hierarchical classification. Unfortunately this measure does not provide a good score to classify those states into communities. We have realized that the fluctuations S_{ij} in $|P_{ik}-P_{jk}|$ indeed play the main role for that classification. Let us take $q=1$ and define

$$\bar{d}_{ij} = \frac{\sum_{k \neq i,j}^{n} |P_{ik} - P_{jk}|}{n-2} \quad (7)$$

as the average "distance" between i and j. The standard deviation between those vertices is given by

$$S_{ij} = \left(\frac{1}{n-3} \sum_{k \neq i,j}^{n} [|P_{ik} - P_{jk}| - \bar{d}(i,j)]^2 \right)^{1/2}. \quad (8)$$

As a matter of fact, this quantity gives a better description of the similarity among the vertices because it is a more sensitive measure than the average value in Eq. (6). The importance of those fluctuations to classify vertices into communities may be surmised saying that we should not ask how far away two vertices are, but who are their neighbors.

Second, we explore the behavior of P_{ij}. A low transition probability to go from state i to j means that state j is less accessible from state i. High transition probabilities among states define classes of easily connected states. This is better

understood in terms of $1/P_{ij}$. Since the elements P_{ij} are not necessarily symmetric, we describe how close vertices i and j are by defining distance as the $\min\{1/P_{ij}, 1/P_{ji}\} = 1/\max\{P_{ij}, P_{ji}\} \equiv 1/P^{\max}_{\{ij\}}$. In other words, the quantity $1/P^{\max}_{\{ij\}}$ establishes how close those vertices are and in this way sets levels of transient classes on $G(V,E)$.

Third, in order to have a well-defined class of states we should expect small total transition probability for leaving it. Let us also introduce the notation $P^{\min}_{\{ij\}} \equiv \min\{P_{ij}, P_{ji}\}$. Thus, a large value of $\Delta_{ij} \equiv P^{\max}_{\{ij\}} - P^{\min}_{\{ij\}}$ is a consequence of a small value for the leaving probability $P^{\min}_{\{ij\}}$ and large value for $P^{\max}_{\{ij\}}$.

Therefore, we achieve the desired hierarchical analysis by defining heuristically a similarity matrix (or "distance matrix") **D** as

$$D_{ij} = S_{ij} \frac{\max\{\Delta_{ij}, P^{\min}_{\{ij\}}\}}{P^{\max}_{\{ij\}}}. \tag{9}$$

Comparative values of $P^{\min}_{\{ij\}}$, for different (i,j) pairs, may be translated as a penalty when they are rather large, which has an intimate connection with Δ_{ij}. Thus, the maximum between Δ_{ij} and $P^{\min}_{\{ij\}}$ enters in Eq. (9) as an extra term to help set a similarity scale. As we will show in the next sections, the symmetric matrix **D** is able to unveil a hierarchical structure of states.

IV. EVALUATING COMMUNITY IDENTIFICATION

To understand the meaning of those transient classes we investigate in some examples the structure of $G(V,E)$ encoded by the similarity matrix. Our analysis reveal well-defined classes of vertices. They occur at different levels of the hierarchical tree under D_{ij} with the interesting interpretation of communities, i.e., with the structure of well-defined subnetworks.

A. Performance on artificial community graphs

Before discussing a particular issue on how to implement the algorithm we report its performance on graphs with a well-known fixed community structure [15]. Our method was tested on large number of graphs with $n=128$ vertices and designed to have four communities each with 32 vertices. Each graph is randomly generated with probability p_{in} to connect vertices in the same community and probability p_{out} to those vertices in different communities. Those probabilities are evaluated in order to make the average degree of each vertex equal to 16. The test amounts to evaluate the fraction of vertices correctly classified as a function of z_{out}, the average number of edges a given vertex must have outside of its own community. Our algorithm classifies correctly vertices into the four communities for small values of z_{out}, decreasing its performance towards $z_{\text{out}}=8$. We have, for example, the fractions 0.99 ± 0.01, 0.95 ± 0.01, 0.81 ± 0.02, 0.57 ± 0.03, respectively, for $z_{\text{out}}=5,6,7$, and 8. The error bar was evaluated over 100 randomly generated graphs. Those results are competitive with the performance obtained by the algorithms compared in Ref. [14]. Moreover, we stress that the proposed method is fully parameter independent. Also, its computational cost is limited by the state-of-art in computing the eigenvalues and eigenvectors of symmetric matrices. In general it amounts to initial $O(n^3)$ operations, with subsequent less expensive iterations $O(n^2)$.

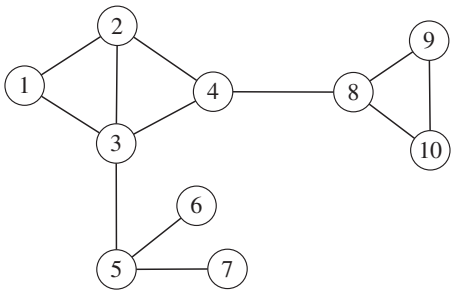

FIG. 1. A simple graph with a treelike subgraph: vertices 5, 6, and 7. Our graphs are drawn using VISONE.

B. A graph with leaves

The method is quite simple and much of the computer time is spent in calculating the eigenvalues and eigenvectors of **L**. All that remains to calculate is the effective resistances in Eq. (4) and, with the elements P_{ij}, the final similarity matrix **D** in Eq. (9). However, some care is needed when the graph presents what we call leaves. This is explained as follows.

We present in Fig. 1 a small graph to display the information contained in the matrix **D** and how to perform the hierarchical analysis. This example shows a graph containing a subgraph with treelike topology. A tree is a connected acyclic graph. In this example, the tree is the subgraph with vertex numbers 5, 6, and 7, which we call leaves. Their effective resistances are $r_{56}=r_{57}=r_{35}=1$ and therefore we have $r_{36}=r_{37}=2$. For treelike subgraphs the effective resistances correspond to the number of edges ℓ_{ij} connecting vertices i and j. Therefore, $r_{ij}=\ell_{ij}$ for acyclic branches. Also $r_{48}=1$ because there is only one way of reaching vertex 8 from vertex 4. On the other hand, whenever we have different paths joining adjacent vertices (i,j), we obtain $r_{ij}<1$ as a consequence of calculating the effective resistance of resistors connected in parallel and in series. For example, $r_{89}=r_{8(10)}=r_{9(10)}=0.6667$. To unveil the hierarchical structure of graphs with leaves, i.e., vertices with degree one, we need to proceed as explained below because well-defined transient classes of states are only identified for graphs with no local treelike topology. Suppose we start with a graph with m vertices $(m=10)$. If the graph has leaves, we collect leaf after leaf to remove acyclic branches and we end up with a reduced number of vertices n $(n<m)$ [32]. After collecting all leaves, we work with the Laplacian matrix of order n obtained from the reduced adjacency matrix. During this process we keep track of the original labels and from where we have removed leaves. The hierarchical structure of this example is shown in Fig. 2 as a dendrogram where we have joined the previously removed leaves (6, 7, and 5) to vertex 3 because they naturally belong to the same community as vertex 3 does. All

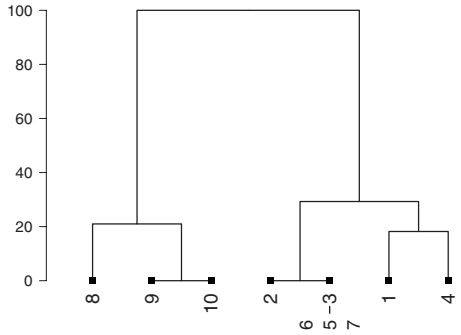

FIG. 2. The community structure of the graph in Fig. 1 is depicted as a hierarchical tree or dendrogram obtained with the complete linkage method based on the similarity matrix **D**. Our dendrograms are drawn with the data plotting package and programming language R.

presented dendrograms have their similarity (y axis) **D** scaled to be in the range (0, 100). This allows a comparative display of their branches with different networks as exemplified in Fig. 6.

C. Zachary karate club network

To illustrate further the meaning of transient classes on $G(V,E)$ from global information carried out by **D** we analyze two well-known networks that have been reported in the literature.

The first example (Fig. 3) corresponds to the network of members of the karate club studied by Zachary [22] and analyzed by different community finding techniques [13–15,17,18,20,21,33]. This network contains a single leaf, member 12. Our analysis leads to the hierarchical structure shown in Fig. 4 by means of a hierarchical clustering tree, defining communities at different levels. The two main communities reproduce exactly the observed splitting of the Zachary club. Interestingly, a smaller community presented by the hierarchical tree is clearly identified in Fig. 3. Its members are displayed with shaded circles. This small group is influenced only by its members and has a direct interaction with the instructor.

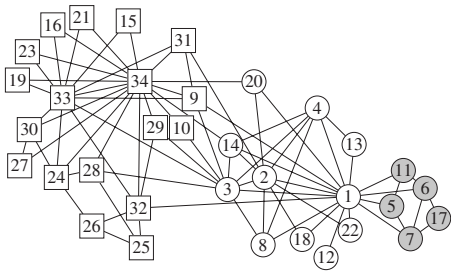

FIG. 3. The karate club network studied by Zachary. Individual numbers represent the members of the club and edges represent their relationships outside the normal activities of the club. Squares and circles indicate the observed final splitting of the karate club into two communities led by the administrator (34) and the instructor (1). A clear further splitting is identified with shaded circles.

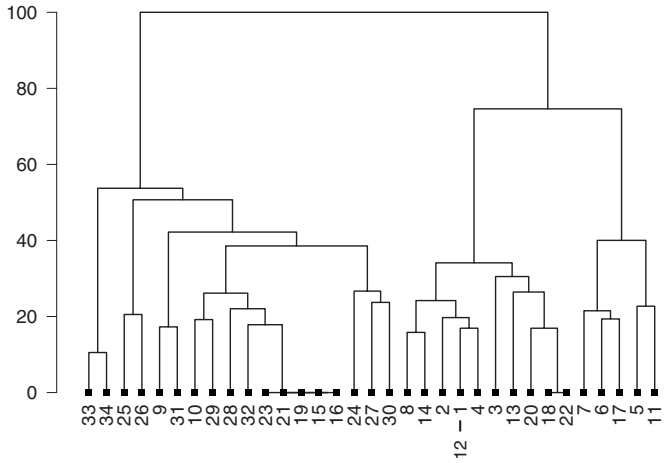

FIG. 4. The hierarchical structure of the network in Fig. 3 is shown as a dendrogram obtained with the complete linkage method. It correctly identifies the two main communities of the karate club.

D. Ravasz and Barabási square hierarchical network

The second example is shown in Fig. 5. It was designed by Ravasz and Barabási [7,23] as a prototype of a hierarchical organization with scale-free topology and high modularity we may encounter in real networks. The network in Fig. 5 was built with the module in (a). A similar network but with more connections between vertices can be built with the module in (b). The study of D_{ij} reveals community structures at different hierarchical levels in Fig. 6, respectively, for the networks generated with the modules (a) and (b).

The hierarchical trees present similar structures, but the hierarchical levels in both figures clearly display different network formation patterns. Moreover, the hierarchical formation pattern of $G(V,E)$ with branches at different heights may be seen as a measure of how cohesive those subgroups are. Thus, the normalized scale for D_{ij} can be used to set levels of cohesiveness related to community formation.

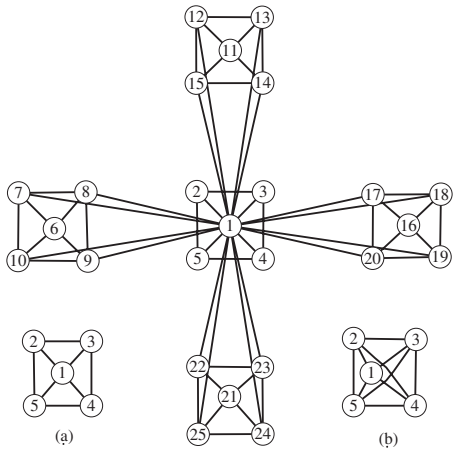

FIG. 5. The deterministic hierarchical scale-free model with $n=5$ vertices proposed by Ravasz *et al.* [7,23]. It is built by generating replicas of the small 5-vertex module (a) shown on the left-hand side.

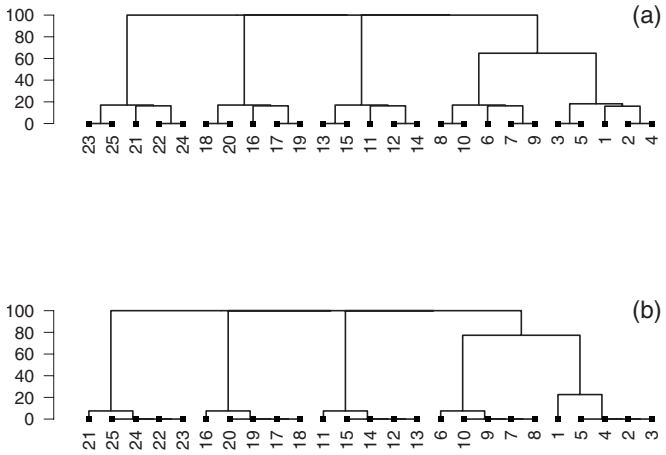

FIG. 6. Hierarchical structures for the formation patterns in Fig. 5. Dendrogram (a) refers to the network generated with module (a) in Fig. 5 whereas dendrogram (b) refers to the network generated with module (b).

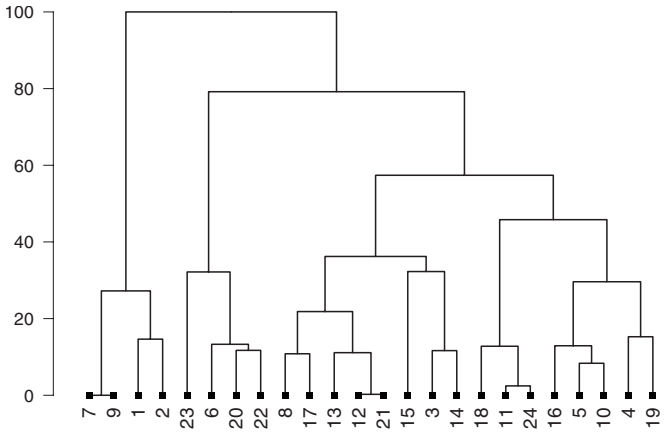

FIG. 7. Hierarchical structure of professional discussions among teachers at "Our Hamilton High."

V. WEIGHTS ON THE EDGES

Our method also applies to graphs such that each edge has a positive real number, the weight of the edge. The structure of the graph is now represented by the corresponding weighted adjacency matrix $\mathbf{W}$. It assigns weight $w_{ij}>0$ if and only if i and j are connected vertices and 0, otherwise. The concept of the Laplacian matrix extends directly to weighted edges, $\mathbf{L}(G)=\mathbf{E}(G)-\mathbf{W}(G)$, where $E_{ii}=\Sigma_{j=1}^{n}w_{ij}$ is the diagonal weighted matrix whose values are the total weight of the edges adjacent to vertex i. Again, $\mathbf{L}(G)$ is a real symmetric matrix where the row sums and the column sums are all zero. Thus, we have the same spectral properties as the particular case of equal weight $w_{ij}=1$ for all adjacent vertices i and j. Therefore, the method presented to unweighted graphs extends naturally to weighted ones with no change in the algorithm.

A. Performance on artificial community weighted graphs

We have also verified the performance of our method on weighted graphs with fixed community structure [34]. Our test is performed on the same artificial graphs randomly generated as described in Sec. IV A. The computer generated graphs have 128 vertices and are divided into four groups of 32 vertices. Here, edges among vertices are randomly chosen such that the average degree is fixed at 16. The test is performed for the most difficult situation where $z_{\text{out}}=z_{\text{in}}=8$. That is, each vertex has as many adjacent connections to inside as to outside its community. For each graph, we attach a weight $w>1$ to the edges inside each community and keep the fixed weight 1 for those edges which lie between communities. We evaluate again the fraction of vertices classified correctly as a function of w. As w increases from the starting value 1, the weights enhance the community structure. This is clearly highlighted by our method. Our performance is demonstrated by the following fractions of correctly classified vertices, 0.89, 0.94, 0.97, and 0.98, respectively, for $w=1.4, 1.6, 1.8$, and 2. The average fractions were calculated over 100 randomly generated graphs, with error bars smaller than 0.01.

B. Identifying cohesive subgroups

As an example, we apply our method to the analysis of weighted interactions related to the engagement of teachers in professional discussions [35]. This is a social network with $n=24$ members. The edges of the network are weighted in accordance with the number of professional discussions between pairs of teachers in a high school, called "Our Hamilton High," during the 1992–1993 school year. Teachers were asked to list and weight the frequency of their discussions in that school with at most five other teachers. This way of attributing weights leads to a directed network. The weights should follow a scale running from 1, for discussions occurring less than once a month, to the largest weight value 4, for almost daily discussions. In this social network, each vertex number contains characteristics of that teacher such as gender, race, subject field, room assignment, among others. To perform our analysis we define the weight of each edge as the average of two values placed on the edges by the two interacting teachers. Thus, the weighted network is characterized by edges with real values in the range 0.5–4 that represent the weighted interactions among the members of the school. The community structure revealed by our analysis is represented by the dendrogram in Fig. 7. Its structure exhibits the formation of several communities. For comparison with the results in Ref. [35], we identify the four main groups. The study of their members reveals largely common race and gender, in accordance with Ref. [35]. However, we differ somewhat in identification of members in each group. This may owe to the fact that our network is weighted by an average process while Frank in Ref. [35] handles it in the original directed form.

VI. CONCLUSIONS

In conclusion, random walks on graphs in connection with electrical networks identify a topological property of $G(V,E)$: transient classes of vertices which we interpret as

communities in the original graph. Here we emphasize that those special classes of vertices are a direct consequence of effective transition probabilities, which display a global perspective about the map of interactions that characterize the entire graph. We demonstrate performance in identifying community structures in some examples which are benchmarks for algorithm validation. Moreover, our method is free of the procedure of tuning a parameter to identify communities. Our criterion to define communities depends only on $G(V,E)$ and not on any explicit definition of what a community structure must be.

It is likely that our proposed algorithm may produce new insights into large graphs. In this case, applications include protein-protein interactions and the compartment identification in food-web structures. The visual information about how members form communities along the hierarchical tree may permit identification and characterization of cohesive communities.

Note added. Recently it was called to our attention about a procedure that reduces the size of networks preserving modularity [32]. The procedure corresponds to replace groups of vertices by a single one and in this way we end up with a smaller network. The simpler groups of vertices to be reduced are hair (which we call leaf) and triangular hair. It is plausible that all algorithms will benefit from the reduction of (peripheral) motifs.

ACKNOWLEDGMENTS

The author acknowledges valuable discussions with O. Kinouchi, and A.S. Martinez, and the support from the Brazilian agencies CNPq (Contract No. 303446/2002-1) and FAPESP (Contract No. 2005/04067-6).

[1] A.-L. Barabási, Nat. Phys. **1**, 68 (2005).
[2] R. Albert and A.-L. Barabási, Rev. Mod. Phys. **74**, 47 (2002).
[3] S. N. Dorogovtsev and J. F. F. Mendes, Adv. Phys. **51**, 1079 (2002).
[4] M. E. J. Newman, SIAM Rev. **45**, 167 (2003).
[5] S. Boccaletti, V. Latora, Y. Moreno, M. Chavez, and D.-U. Hwang, Phys. Rep. **424**, 175 (2006).
[6] P. Holme, M. Huss, and H. Jeong, Bioinformatics **19**, 532 (2003).
[7] E. Ravasz, A. L. Somera, D. A. Mongru, Z. N. Oltvai, and A.-L. Barabási, Science **297**, 1551 (2002).
[8] R. Guimerà and L. A. N. Amaral, Nature (London) **433**, 895 (2005).
[9] M. Vendruscolo, N. V. Dokholyan, E. Paci, and M. Karplus, Phys. Rev. E **65**, 061910 (2002).
[10] L. H. Greene and V. A. Higman, J. Mol. Biol. **334**, 781 (2003).
[11] A. R. Atilgan, P. Akan, and C. Baysal, Biophys. J. **86**, 85 (2004).
[12] N. A. Alves and A. S. Martinez, Physica A **375**, 336 (2007).
[13] M. E. J. Newman, Eur. Phys. J. B **38**, 321 (2004).
[14] L. Danon, A. Diaz-Guilera, J. Duch, and A. Arenas, J. Stat. Mech.: Theory Exp. (2005) P09008.
[15] M. E. J. Newman and M. Girvan, Phys. Rev. E **69**, 026113 (2004).
[16] S. Van Dongen, Ph.D. thesis, University of Utrecht, The Netherlands, 2000; D. Gfeller, P. De Los Rios, A. Caflisch, and F. Rao, Proc. Natl. Acad. Sci. U.S.A. **104**, 1817 (2007).
[17] H. Zhou, Phys. Rev. E **67**, 041908 (2003); **67**, 061901 (2003).
[18] F. Wu and B. A. Huberman, Eur. Phys. J. B **38**, 331 (2004).
[19] M. E. J. Newman, Phys. Rev. E **74**, 036104 (2006).
[20] L. Donetti and M. A. Muñoz, J. Stat. Mech.: Theory Exp. (2005) P10012.
[21] F. Radicchi, C. Castellano, F. Cecconi, V. Loreto, and D. Parisi, Proc. Natl. Acad. Sci. U.S.A. **101**, 2658 (2004); C. Castellano, F. Cecconi, V. Loreto, D. Parisi, and F. Radicchi, Eur. Phys. J. B **38**, 311 (2004);
[22] W. W. Zachary, J. Anthropol. Res. **33**, 452 (1977).
[23] E. Ravasz and A.-L. Barabási, Phys. Rev. E **67**, 026112 (2003).
[24] R. Grone, Linear Algebr. Appl. **150**, 167 (1991), and references therein.
[25] S. L. Campbell and C. D. Meyer, *Generalized Inverses of Linear Transformations* (Dover, New York, 1991).
[26] A. Baltz and L. Kliemann, Lect. Notes Comput. Sci. **3418**, 373 (2005).
[27] K. M. Hall, Manage. Sci. **17**, 219 (1970).
[28] B. Bollobás, *Modern Graph Theory* (Springer-Verlag, New York, 1998).
[29] P. G. Doyle and J. L. Snell, http://math.dartmouth.edu/ doyle/docs/walks/walks.ps
[30] W. Xiao and I. Gutman, Theor. Chem. Acc. **110**, 284 (2003).
[31] I. Gutman and W. Xiao, Cl. Sci. Math. Nat, **129**, 15 (2004).
[32] A. Arenas, J. Duch, A. Fernández, and S. Gómes, e-print arXiv:physics/0702015.
[33] M. E. J. Newman, Phys. Rev. E **69**, 066133 (2004).
[34] M. E. J. Newman, Phys. Rev. E **70**, 056131 (2004).
[35] K. A. Frank, Soc. Networks **18**, 93 (1996).